

THz makes X-Ray

János Hebling

hebling@fizika.ttk.pte.hu

Gábor Almási, József Fülöp, Mechler Mechler, László Pálfalvi,

Zoltán Ollmann, Zoltán Tibai, György Tóth

University of Pécs, Institute of Physics

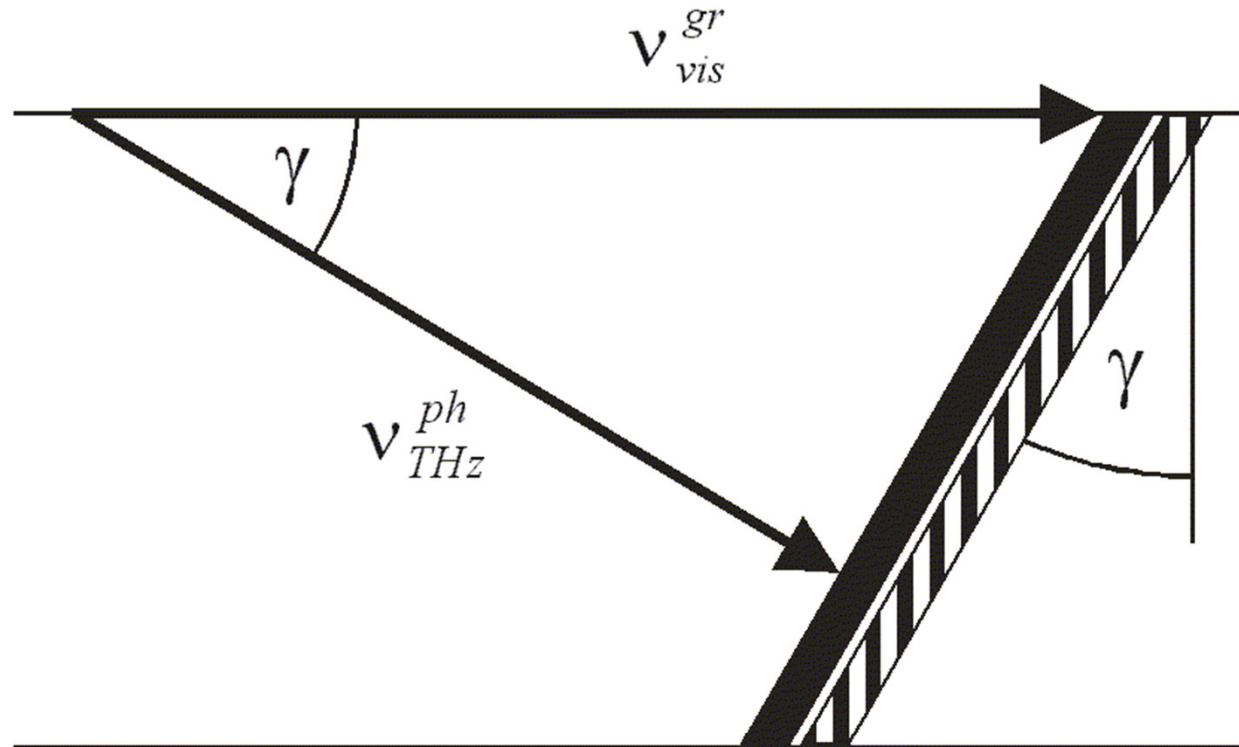


THz Workshop, Argonne National
Laboratory, July 30-31, 2012

Motivation

Velocity matching by tilting of the pump-pulse-front

J. Hebling et al: Optics Exp. **10**, 1161 (2002)



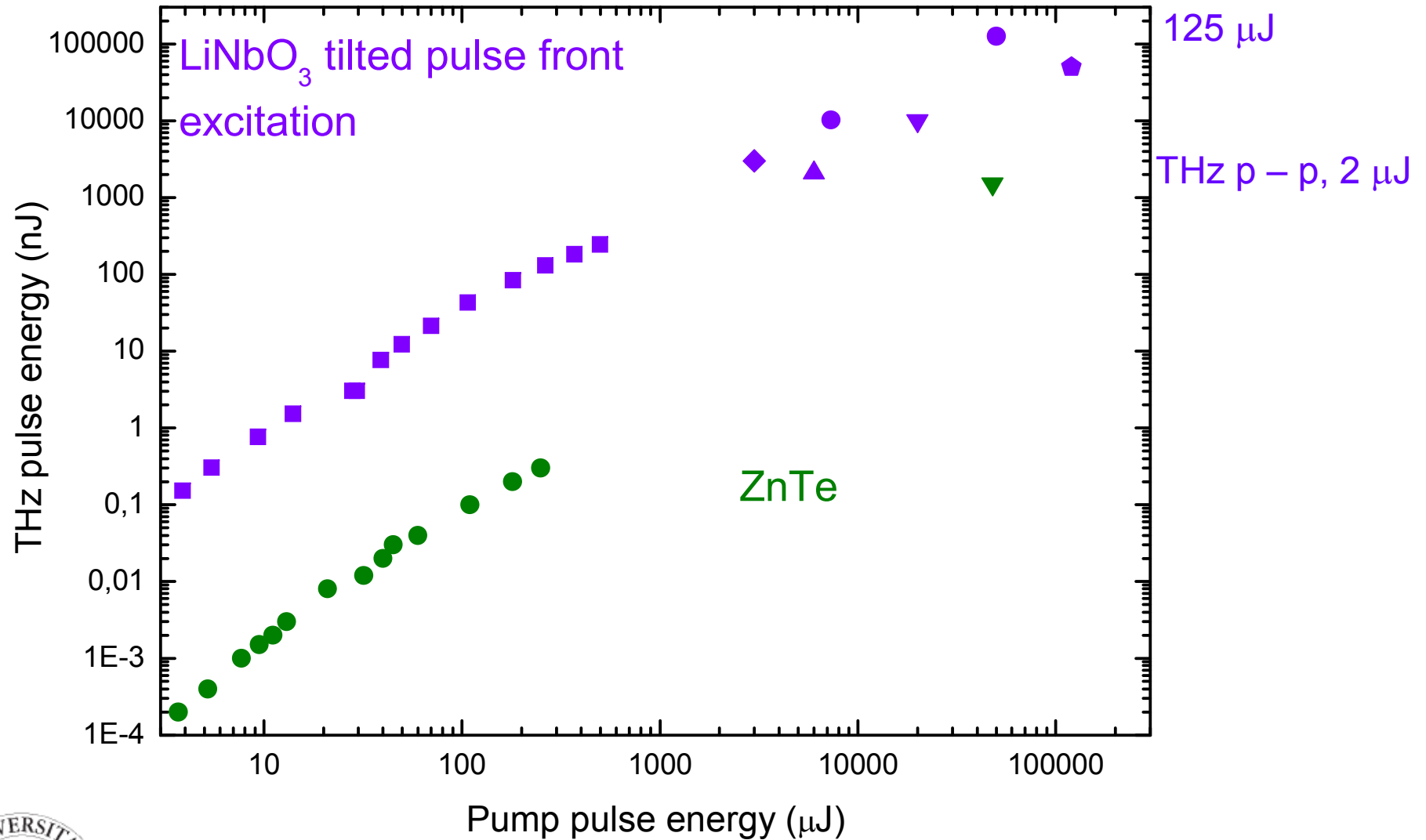
THz radiation excited along the tilted pulse front propagates perpendicularly to this front → velocity matching condition:

$$v_{vis}^{gr} \cdot \cos \gamma = v_{THz}^{ph}$$



Motivation

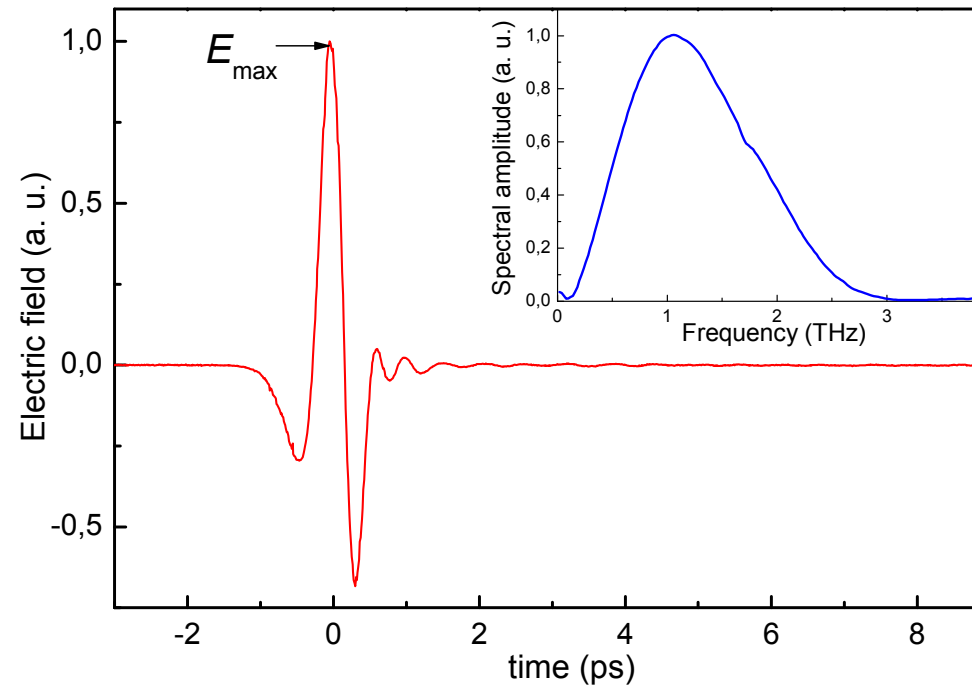
Comparison of THz pulse generation by ZnTe and by LN



Field strength of focused pulses: 0.2 – 1.2 MV/cm

Motivation

Classification of THz pulses by peak electric field (energy)



- Linear (TDTs) THz spectroscopy ($E_{\max} \approx 100 \text{ V/cm} \rightarrow 10 \text{ fJ energia}$)
- High field THz pulses ($E_{\max} \approx 100 \text{ kV/cm} \rightarrow \mu\text{J energia}$)



Applications of high energy THz pulses

- Nonlinear THz optics (self-phase-modulation, Kerr-effect)
- Nonlinear THz spectroscopy (THz pump - probe)
- Electron wave packet sampling (streaking)

Frühling: Nature Photonics **3**, 523 (2009)

Schütte et al.: Opt. Express **19**, 18833 (2011)

- Alignment/orientation of molecules

Stapelfeldt, Seideman: Rev. Mod. Phys. **75**, 543 (2003)

Fleischer et al.: Phys. Rev. Lett. **107**, 163603 (2011)

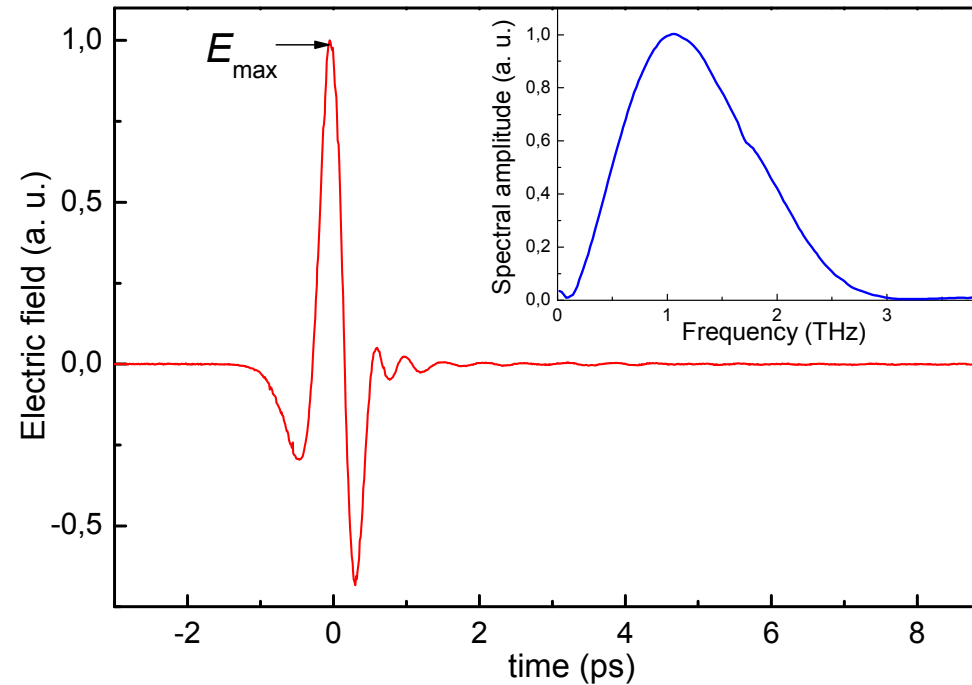
Kitano et al.: PRA **84**, 053408 (2011)

- „Single-shot” THz imaging
- „Single-shot” THz multi spectral imaging



Motivation

Classification of THz pulses by peak electric field (energy)



- Linear (TDTTS) THz spectroscopy ($E_{\max} \approx 100 \text{ V/cm} \rightarrow 10 \text{ fJ energia}$)
- High field THz pulses ($E_{\max} \approx 100 \text{ kV/cm} \rightarrow \mu\text{J energia}$)
- Extreme high field THz pulses
($E_{\max} \approx 100 \text{ MV/cm} \rightarrow 10 \text{ mJ pulse energy}$)



Tomorrow: Alan Fischer: 500 μJ , József Fülöp: 125 μJ

Application possibilities of THz pulses having extreme high electric field strength

- Increasing the cut-off frequency of HHG ($E_{\text{THz}} = 20 \div 80 \text{ MV/cm}$)
- Quasi-phase-matched attosecond pulse generation ($E_{\text{THz}} = 2 \div 6 \text{ MV/cm}$)
- Temporal compression of electron bunches by THz pulses
- Single-cycle (or shaped) pulse generation on the 50 nm – 400 nm wavelength range by coherent Thomson-scattering
- Ultrashort X-ray pulse generation by incoherent Thomson-scattering
- THz driven undulator for UV – X-ray generation



Increasing the cut-off frequency of HHG

HHG in the presence of THz electric field

- Combined THz + IR fields

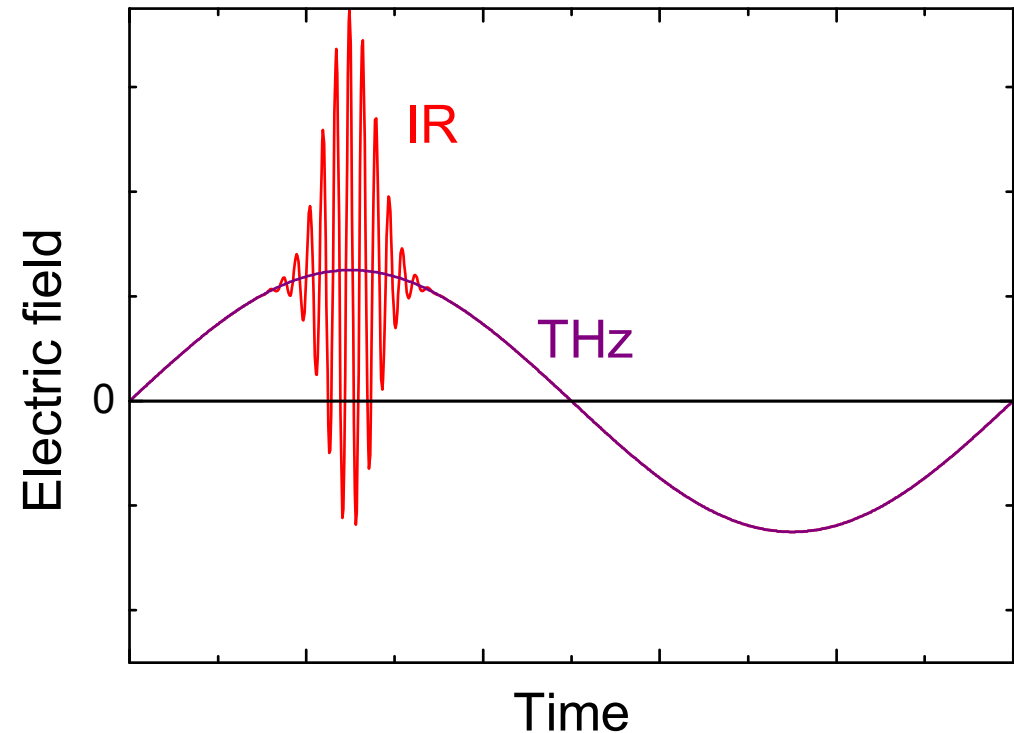
$$E(t) = E_0 \cos(\omega_{IR}t) + E_1 \cos(\omega_{THz}t)$$

- THz broken the symmetry of IR field

Hong et al., Opt. Expr. 2009

- One as pulse per IR cycle
- Spectrum consists both odd and even harmonics

Lewenstein, PRA 1994



E. Balogh et al.: Phys. Rev. A (2011), University of Szeged



Increasing the cut-off frequency of HHG

1560 nm

800 nm

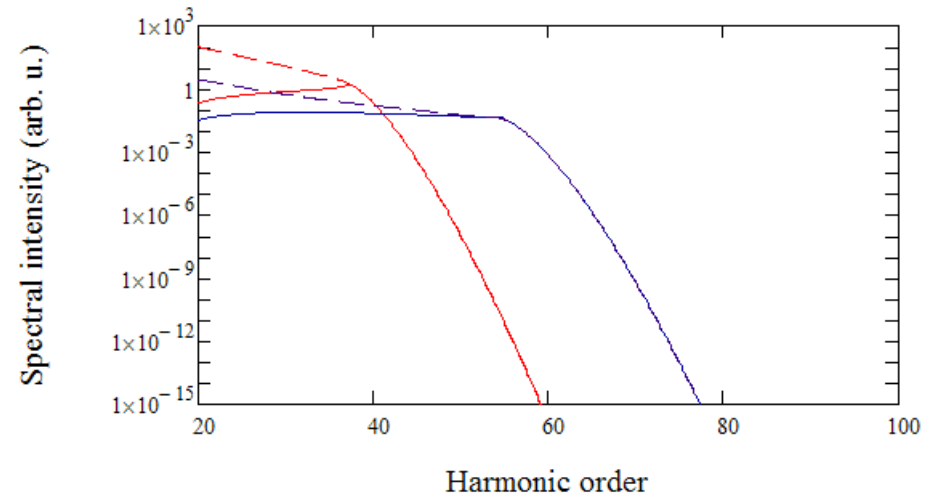
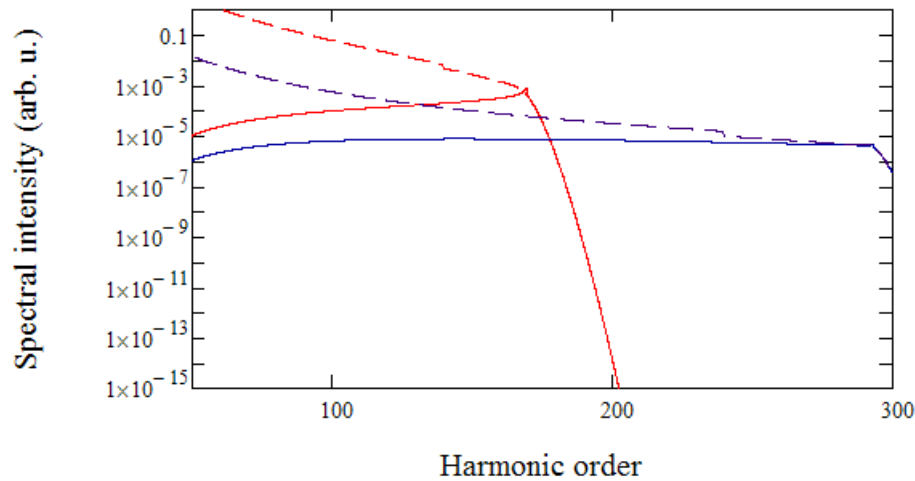
Both odd and even harmonics →→ only one as pulse per IR period

$$I_{\text{IR}} = 2 \times 10^{14} \text{ W cm}^{-2}$$

$$E_1 = 4 \times 10^9 \text{ V/m}$$

$$m_{(\text{coeffindex}_1)} = 37.92 \quad E_{\text{IR}} = 388 \text{ MV/cm}$$

$$E_0 = 3.881 \times 10^{10} \text{ V/m} \quad E_1 = 4 \times 10^9 \text{ V/m}$$



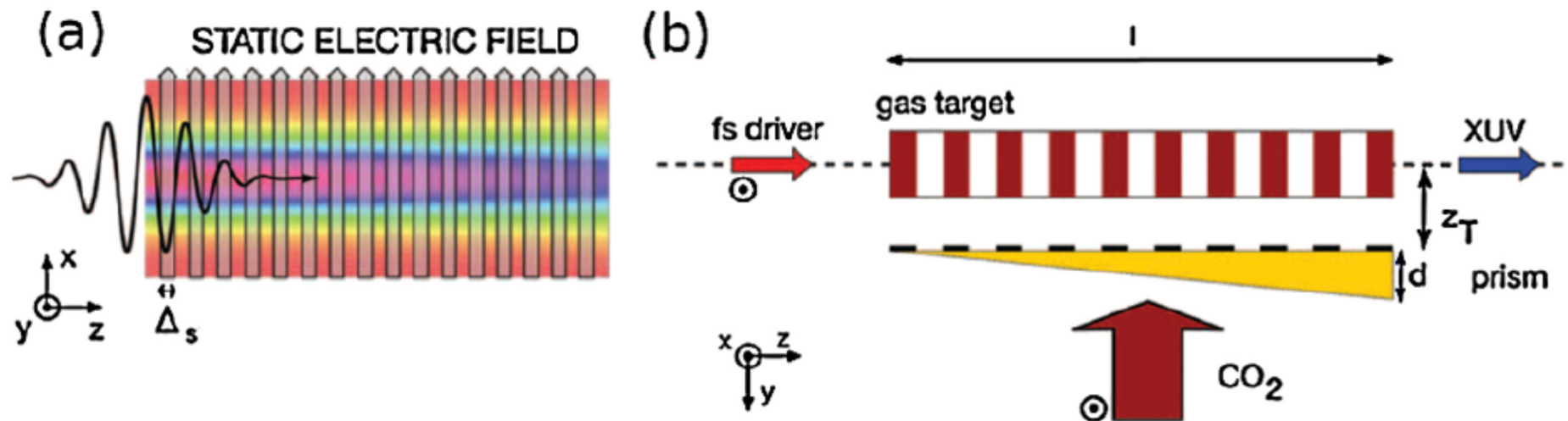
$$E_{\text{THz}} = 0 \dots 40 \text{ MV/cm}$$

E. Balogh et al.: Phys. Rev. A (2011), University of Szeged



Quasi-phase-matched attosecond pulse generation

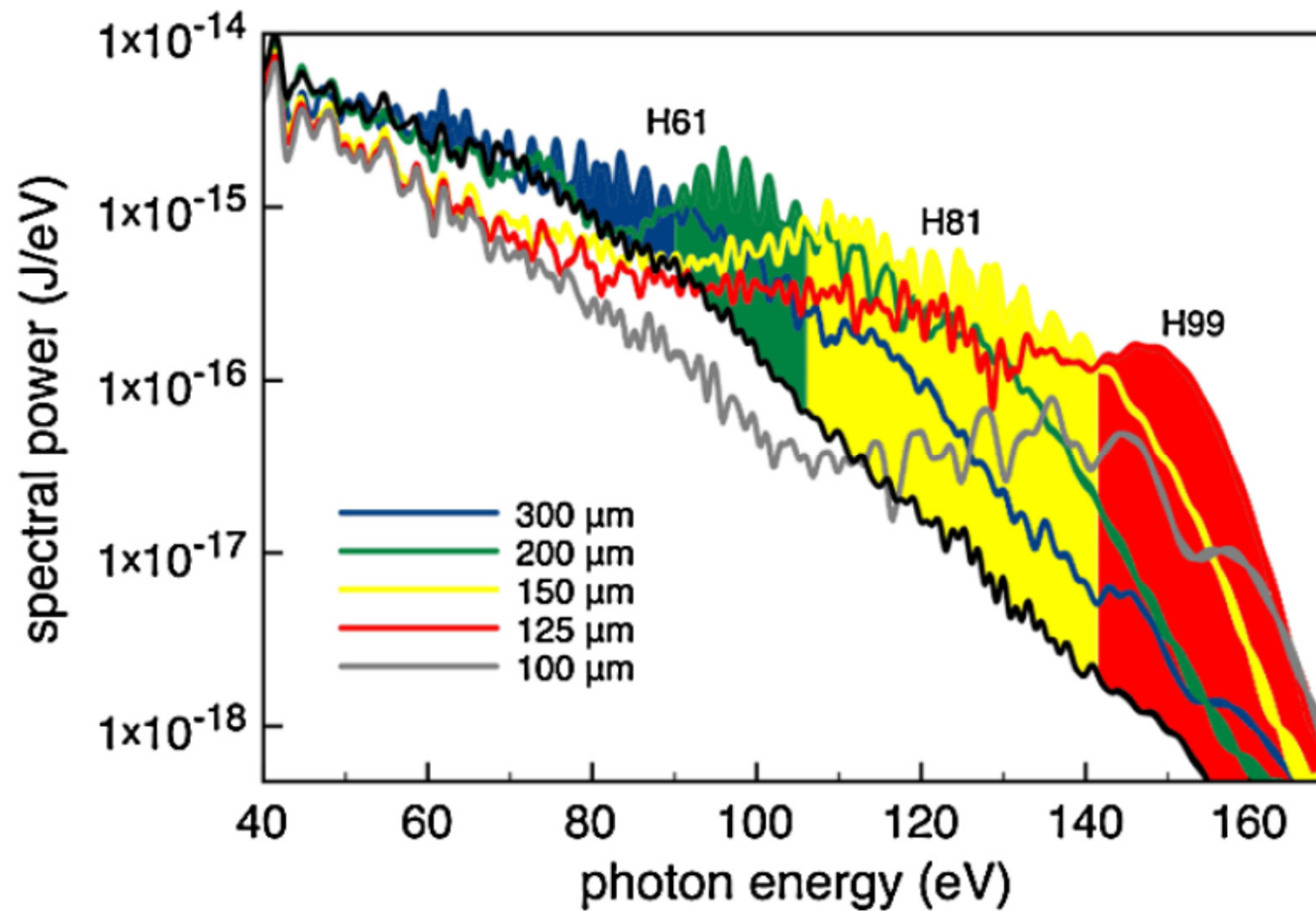
Spatially periodically modulated mid-IR (CO_2 laser)



C. Serrat, J. Biegert: Phys. Rev. Lett. **104**, 073901 (2010)



Quasi-phase-matched attosecond pulse generation

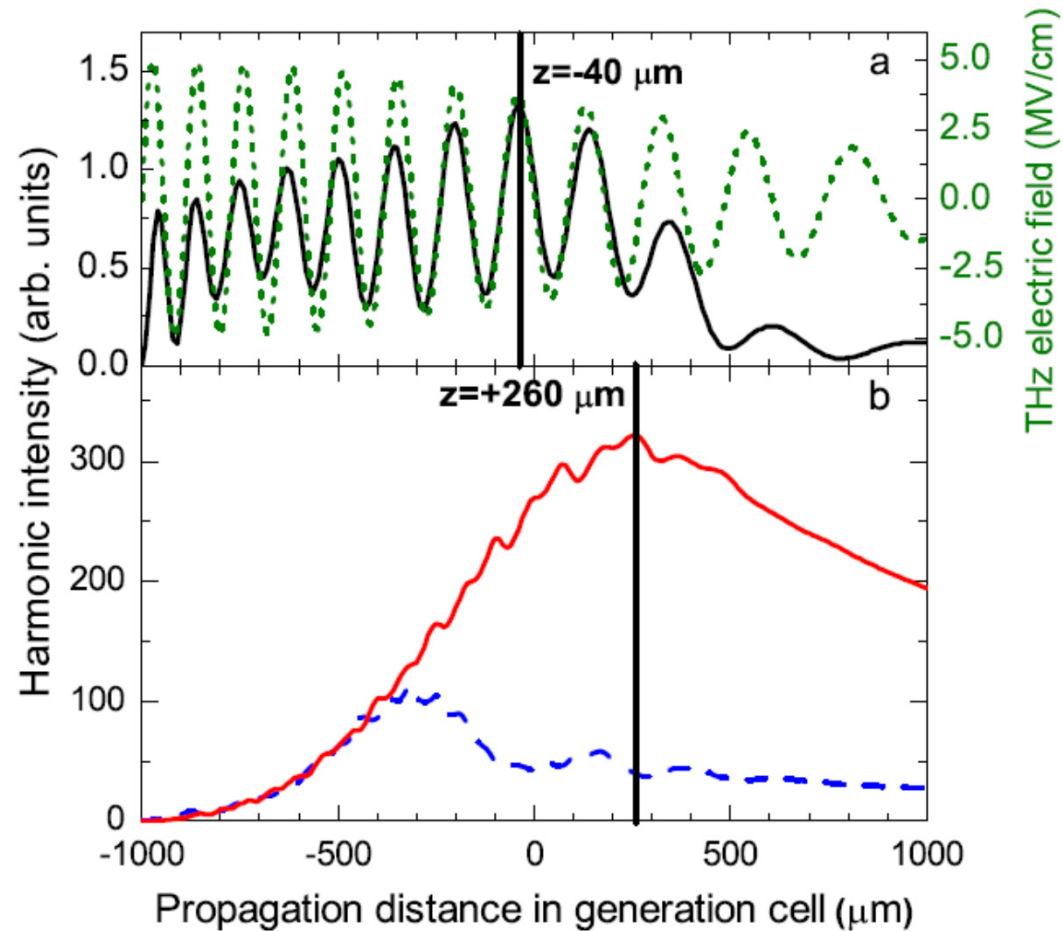


5 fs, 800 nm, Neon, 20 mbar, L=2mm



Quasi-phase-matched attosecond pulse generation

Spatial modulation by (chirped) THz pulse

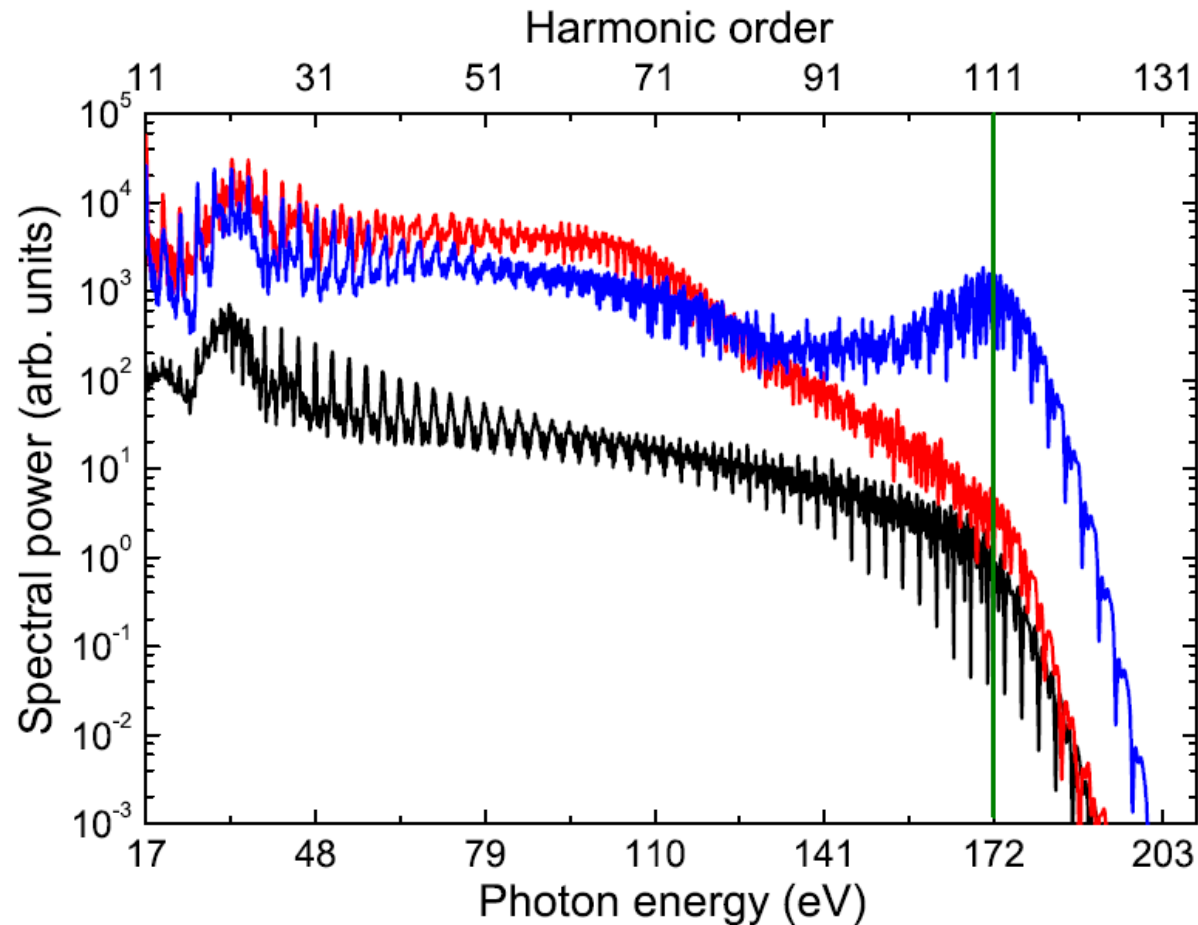


K. Kovács et al.: Phys. Rev. Lett. **108**, 193903 (2012)



Quasi-phase-matched attosecond pulse generation

Spatial modulation by (chirped) THz pulse

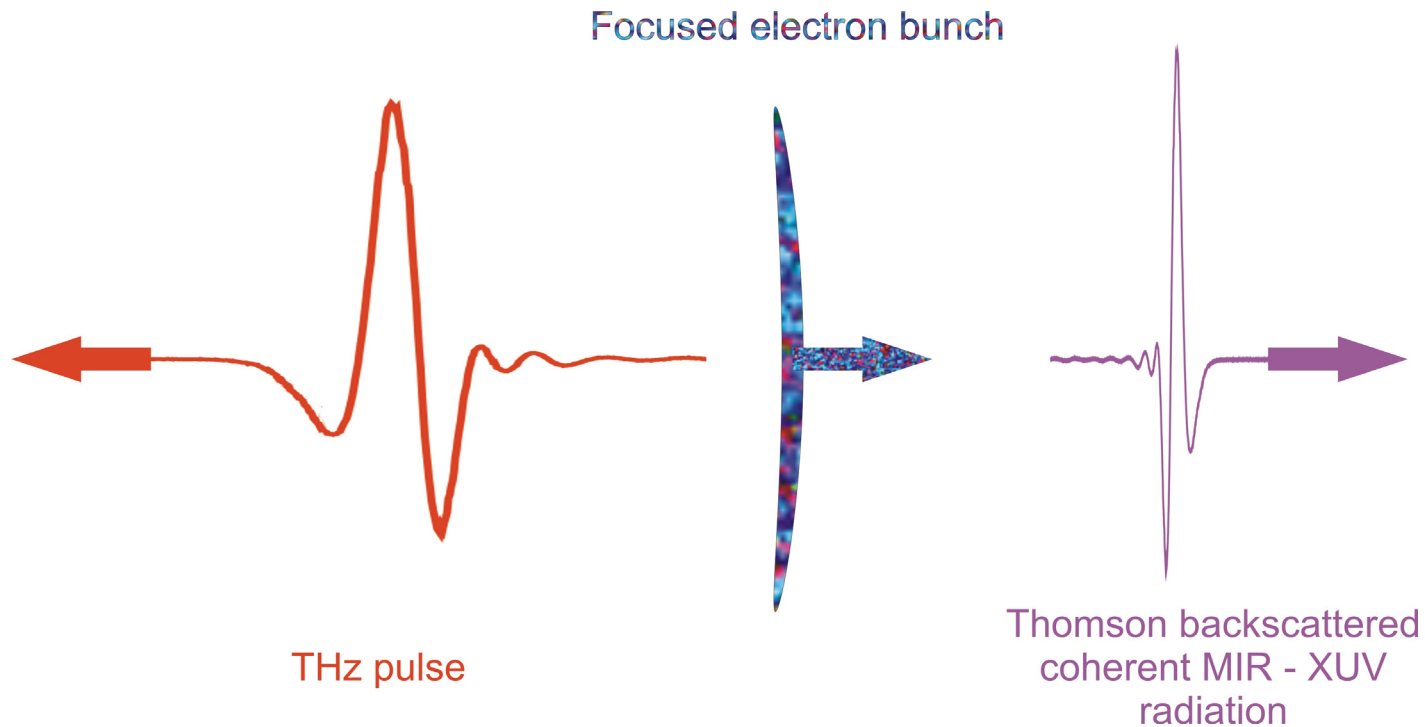


250 times enhancement around the cut-off frequency

$$E_{\text{THz}} = 2 \div 6 \text{ MV/cm}, W_{\text{THz}} = 2 \div 20 \text{ mJ}$$



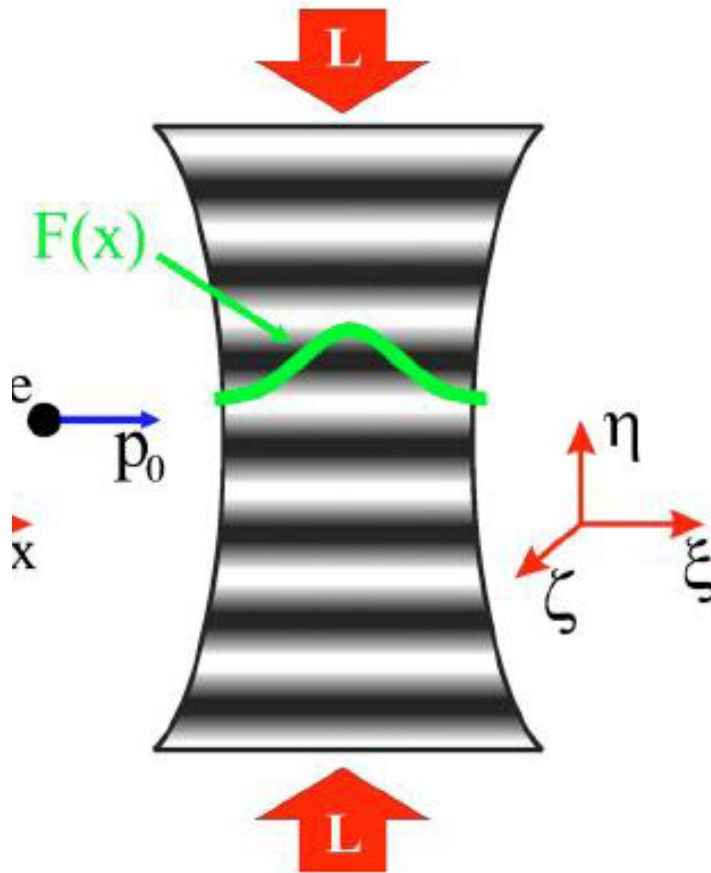
Single-cycle MIR – EUV pulse generation by „coherent” Thomson-scattering



Single-cycle MIR – EUV pulse generation by „coherent” Thomson-scattering



Temporal compression of electron bunches

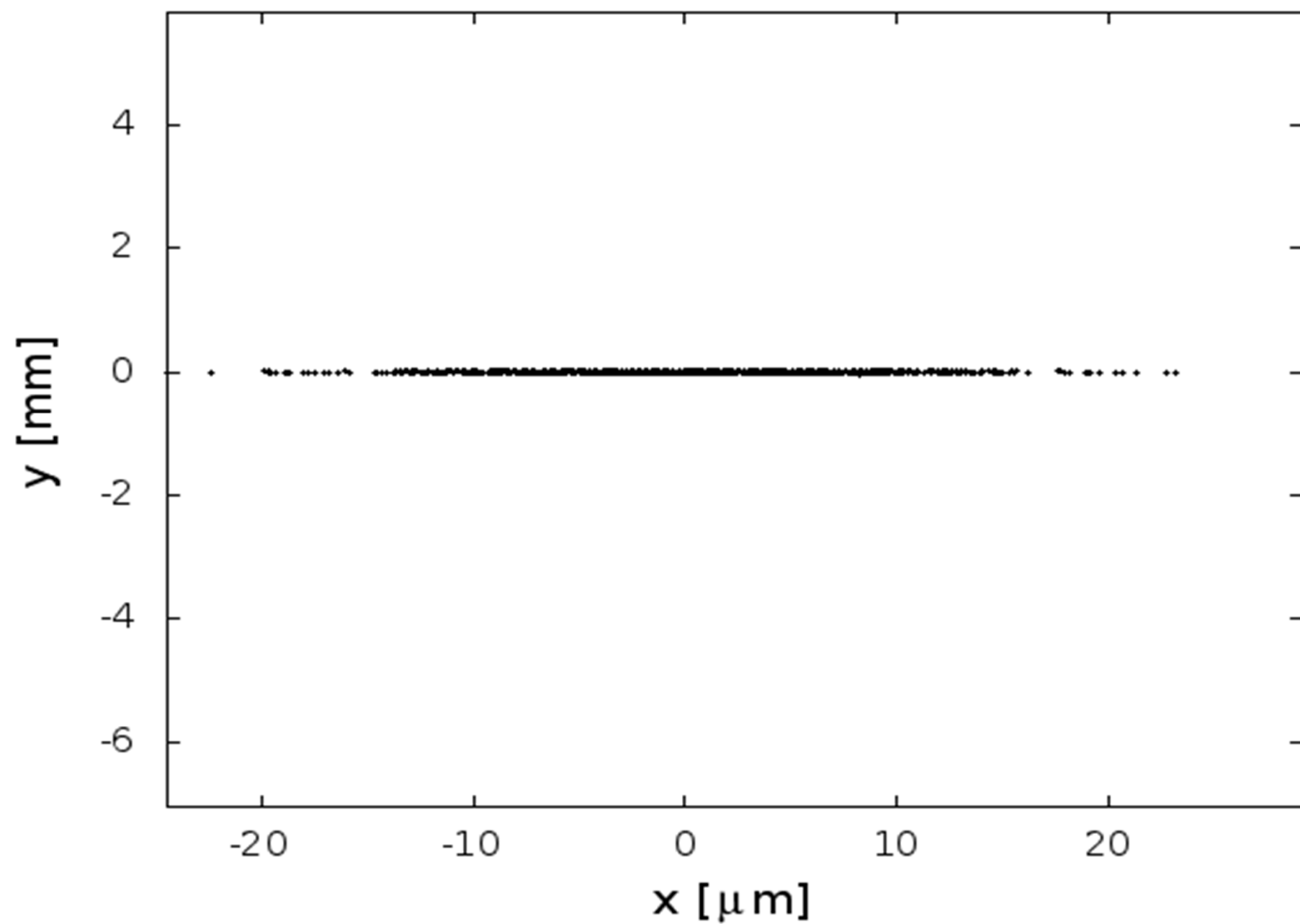


$$\lambda = 1 \mu\text{m}, \Delta t_f = 30 \text{ zS}$$

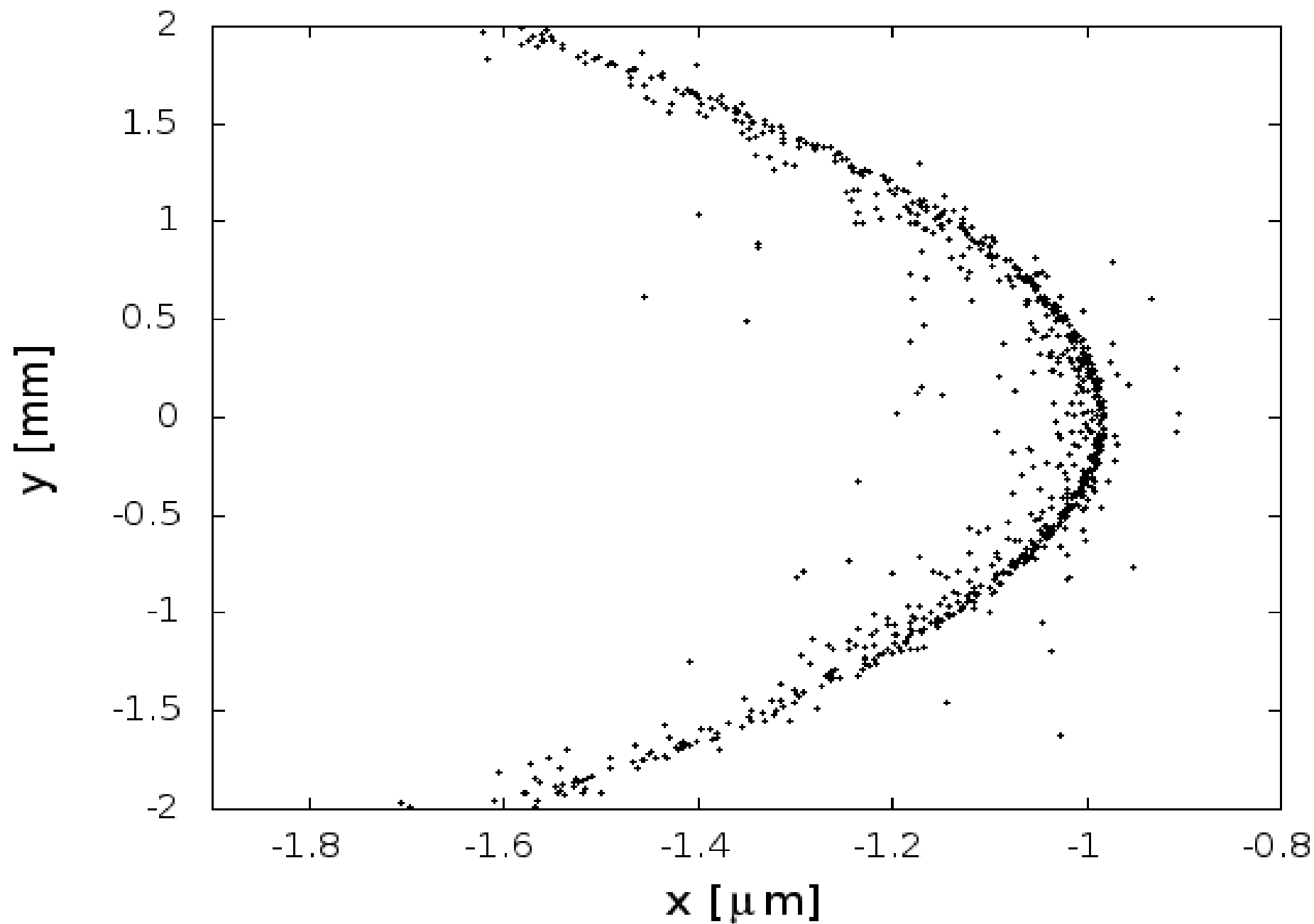
A. E. Kaplan, A. L. Pokrovsky: Optics Express **17**, 6194 (2009)



$\dot{u}t=0.0\text{m}$, $\gamma_{\min}=19.93$, $\gamma_{\max}=20.059$



$\dot{u}t=3.248m$, $\gamma_{\min}=19.93$, $\gamma_{\max}=20.059$



Creation of ultrashort electron bunches by inverse-FEL

Simulation by GPT

parameters of e-bunch:

J. F. Yang: *Jap. J. Appl. Phys.* **44**, 12 (2005)

$\gamma = 62$, $\sigma_E/E = 0.04\%$,

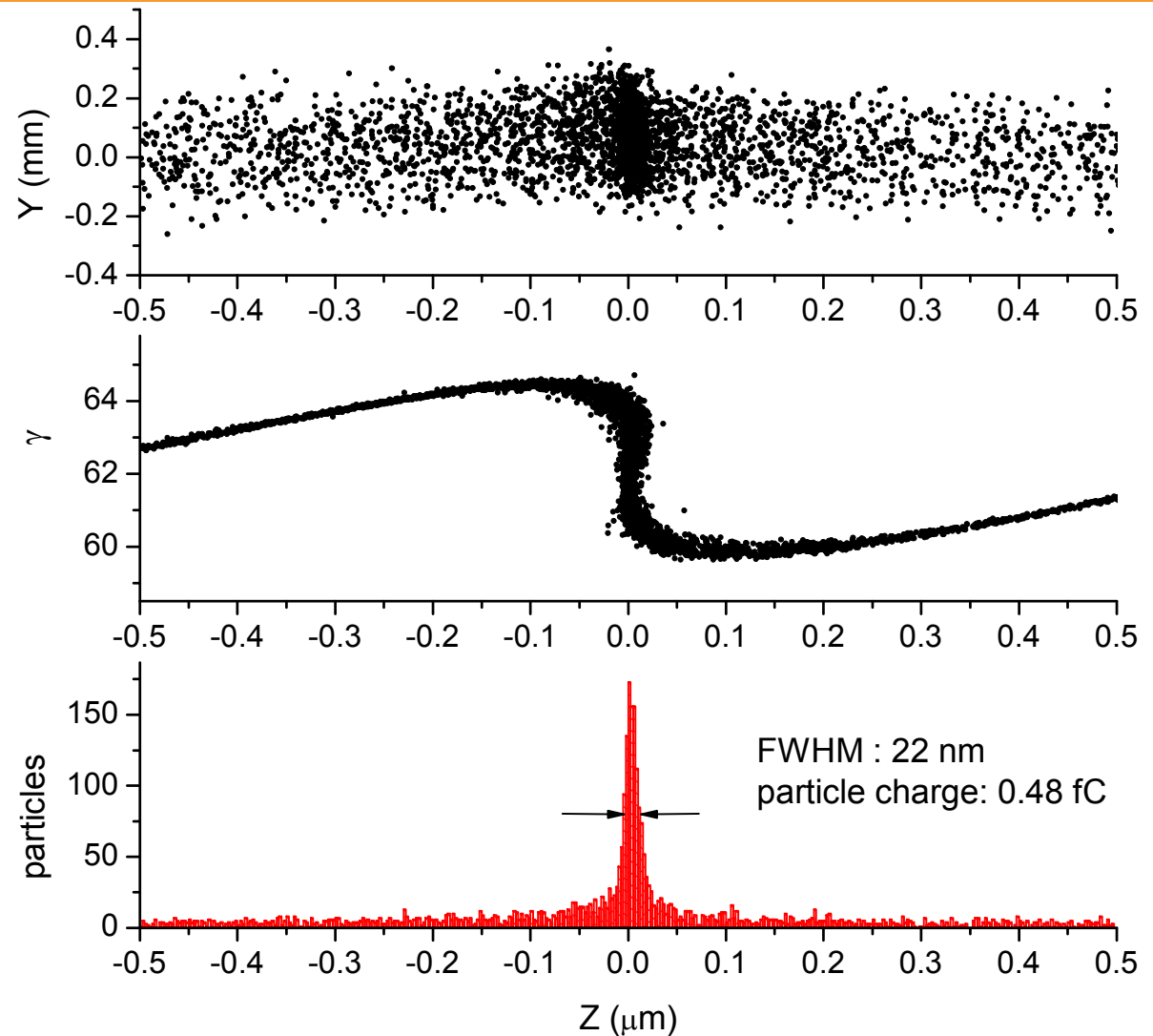
$Q_b = 1$ nC, $\Delta t_b = 1.8$ ps,

$\varepsilon_n = 3.2$ mm

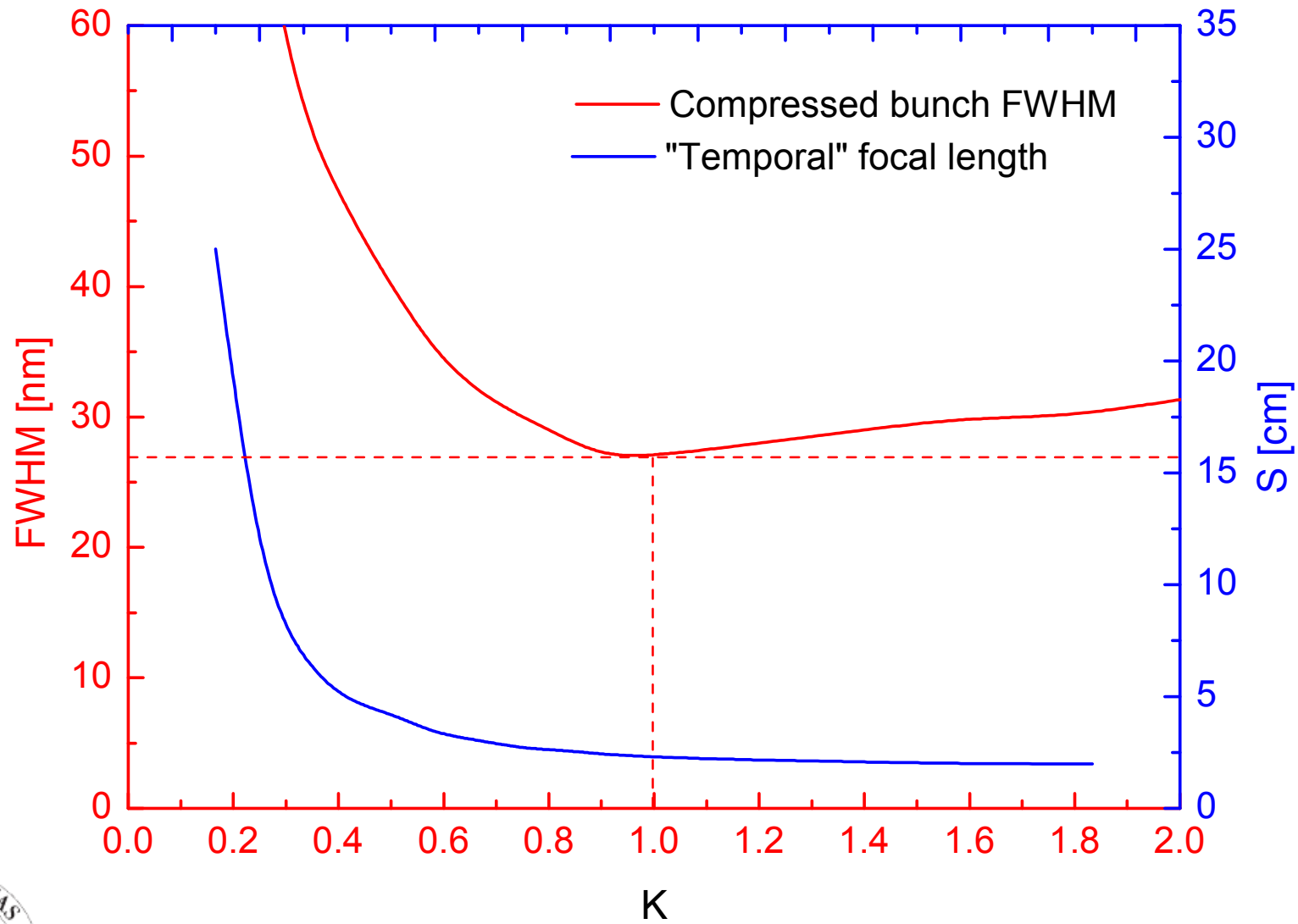
wiggler: one-period, $K = 1$

laser:

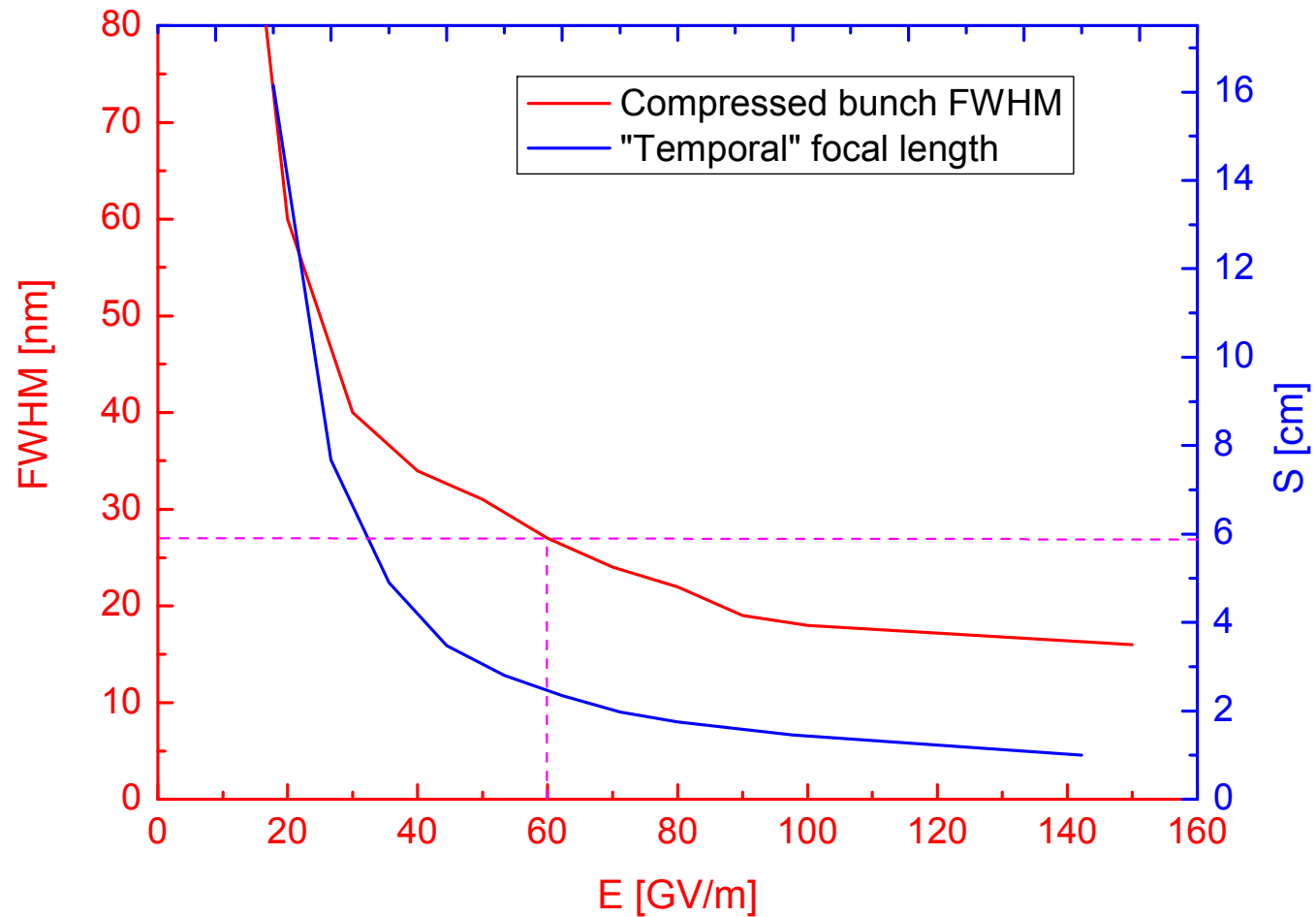
$\lambda = 1.3$ μm , $P = 3.9$ TW



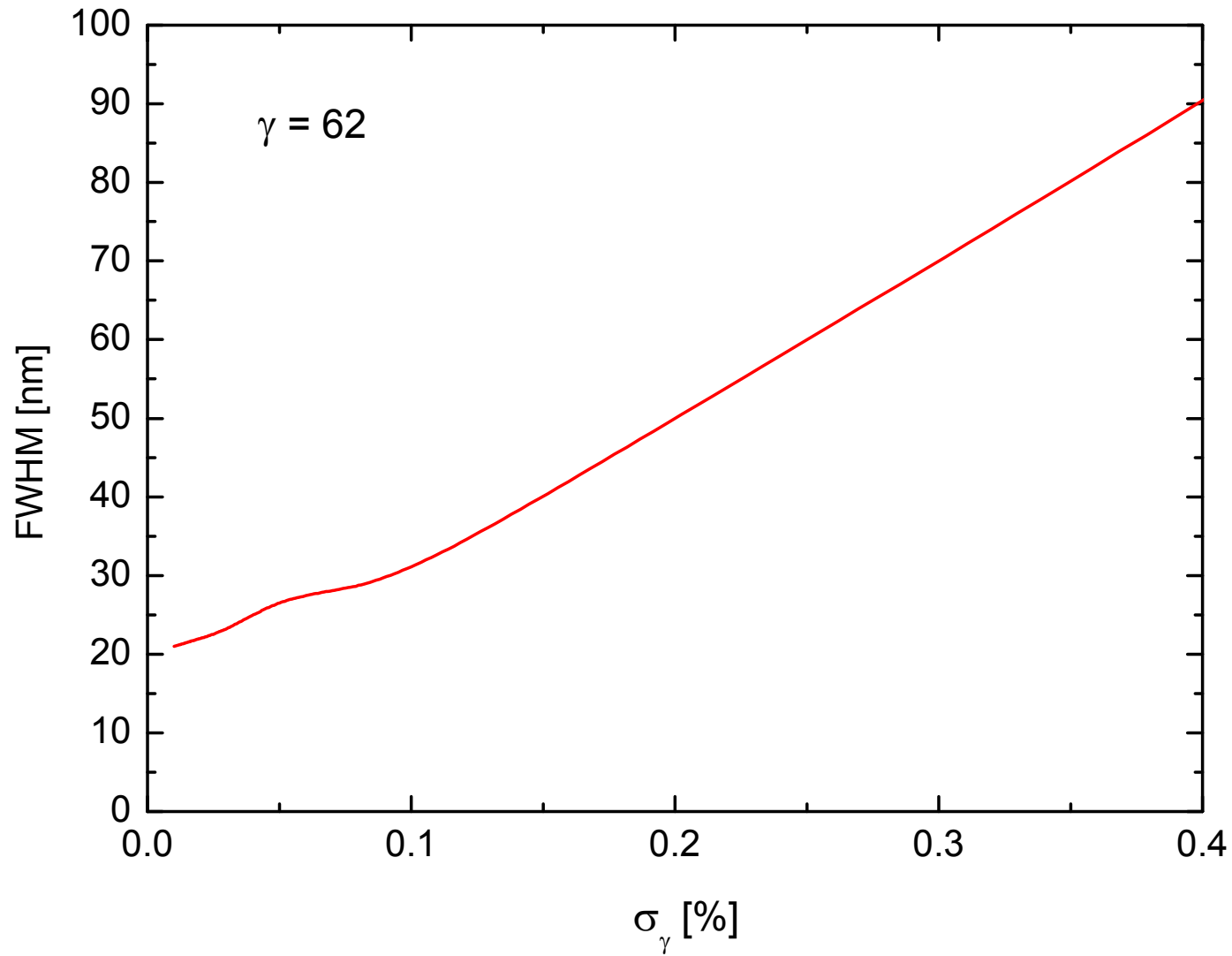
Creation of ultrashort electron bunches by inverse-FEL



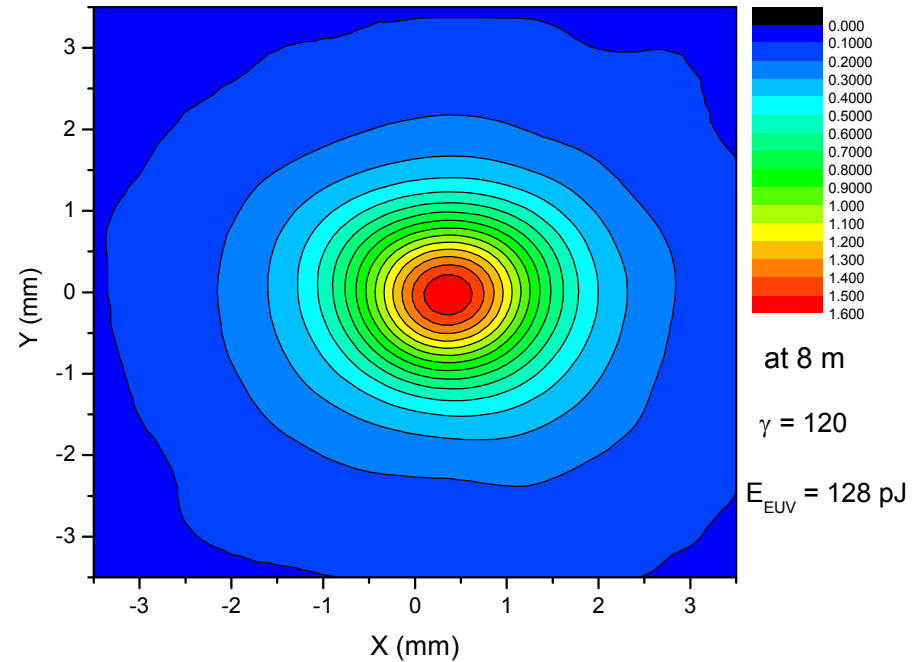
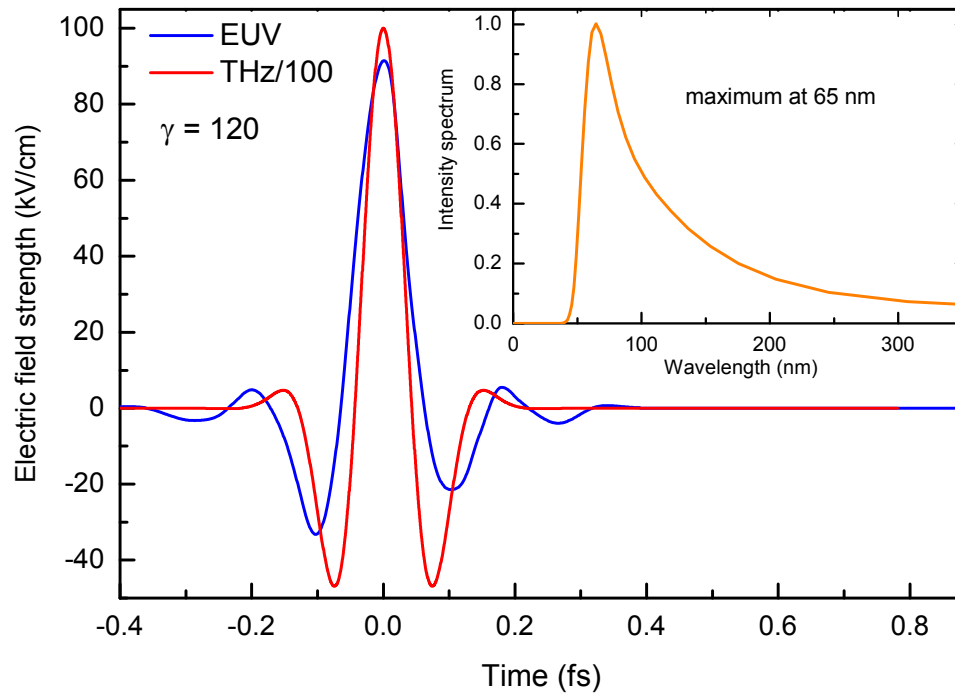
Creation of ultrashort electron bunches by inverse-FEL



Creation of ultrashort electron bunches by inverse-FEL



Single-cycle EUV pulse generation by coherent Thomson-scattering



$$\frac{d\vec{p}}{dt} = q\vec{E} + q\vec{v} \times \vec{B}$$

$$\vec{E} = \frac{q}{4\pi\epsilon_0} (1-\beta) \frac{\vec{R} - R\vec{\beta}}{(R - \vec{R}\vec{\beta})^3} + \frac{q\mu_0}{4\pi} \frac{\vec{R} \times ((\vec{R} - R\vec{\beta}) \times \dot{\vec{\beta}})}{(R - \vec{R}\vec{\beta})^3}$$

$$E_{THZ} = 10 \text{ MV/cm} \quad E_{THZ} = 30 \text{ MV/cm}$$

$$\nu_{THz} = 0.1 \text{ THz} \quad \nu_{THz} = 0.3 \text{ THz}$$

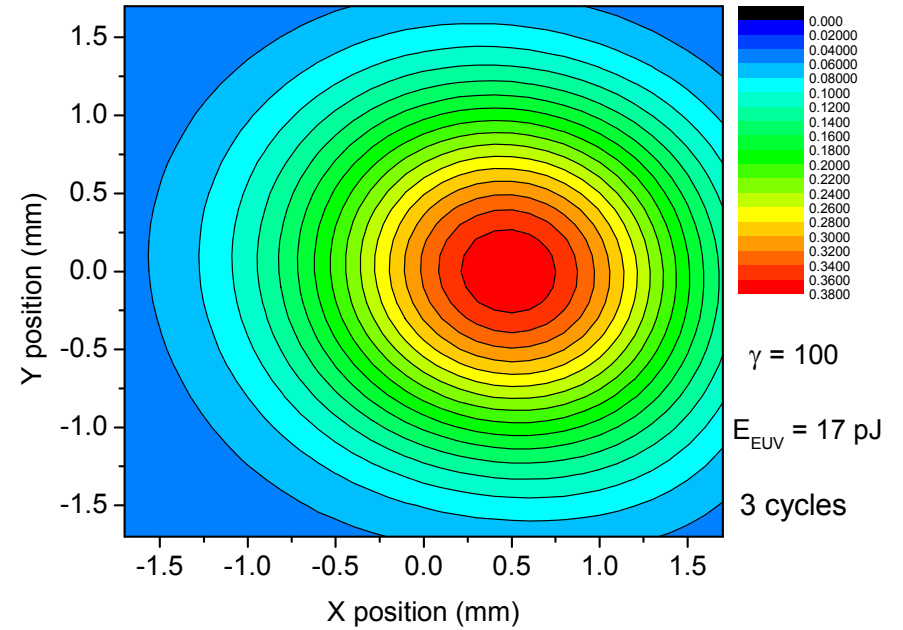
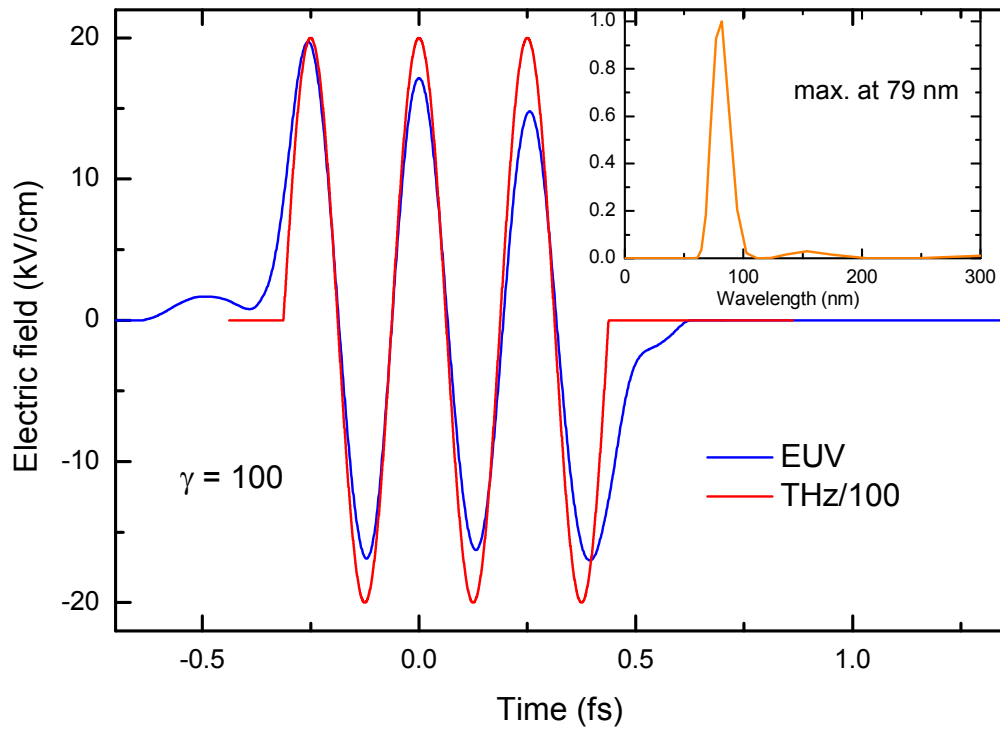
$$\gamma = 120 \quad \gamma = 70$$

$$E_{EUV} = 128 \text{ pJ} \quad E_{EUV} = 15 \text{ pJ}$$

$$\lambda_{EUV} = 65 \text{ nm} \quad \lambda_{EUV} = 85 \text{ nm}$$



Three-cycle EUV pulse generation by coherent Thomson-scattering



3 Å X-ray source by incoherent Thomson-scattering

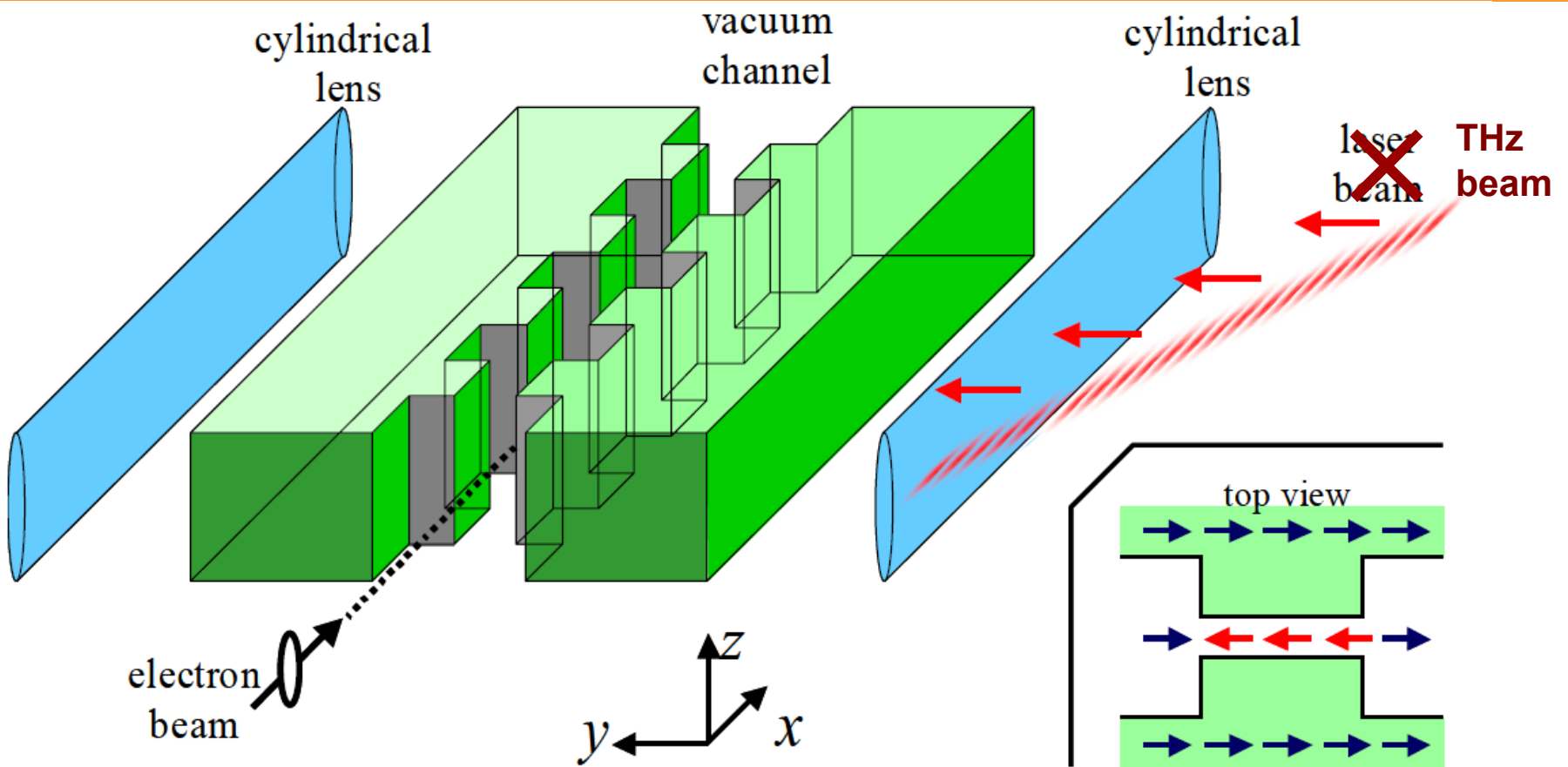
Electron bunch: $Q = 50$ pC, $\gamma = 800$, $\Delta t = 20$ fs [Leemans et al.: N. Phys. 2, 696 \(2006\)](#)

THz pulse: $E = 6$ MV/cm, $\lambda = 750$ μm , $N = 20$, $\Delta t = 12.5$ ps, $W = 24$ mJ

X-ray: $\lambda = 3$ Å, $\Delta t = 20$ fs, $W = 5$ nJ



Electron manipulation by THz pulses



Plettner et al.: Phys. Rev. Spec. Top. – Accel. and Beams **9**, 111301 (2006) acceleration
11, 030704 (2008) undulator
12, 101302 (2009) deflection, focusing



1 GV/m = 10 MV/cm peak field strength is needed!

Laser-plasma accelerator

Summary

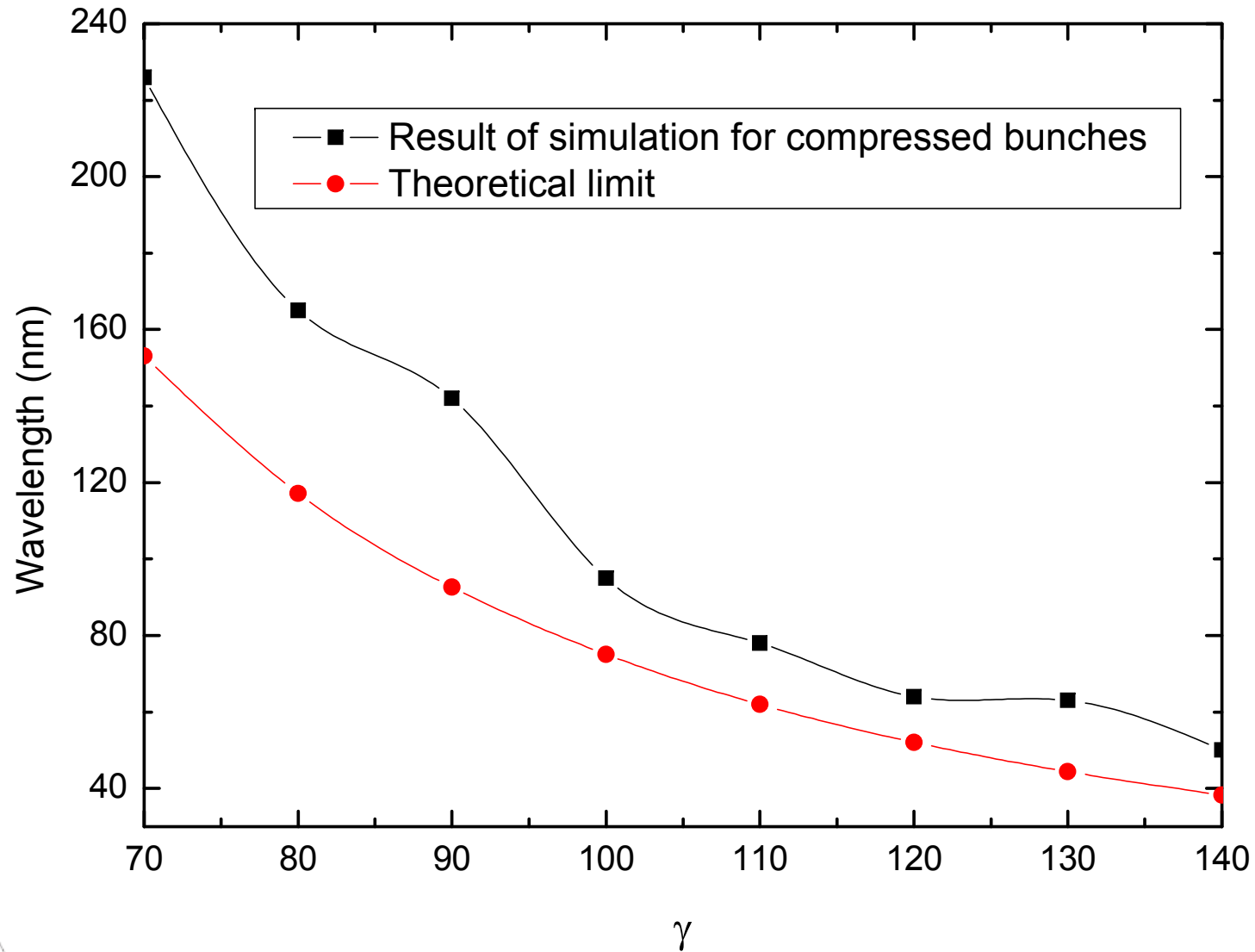
According to simple estimation and numerical simulations, -because of their appropriate wavelength and field strength-, extreme intense THz pulses are suitable for:

- improving attosecond pulse generation
- generation single-cycle or prescribed shaped EUV – VIS pulse generation
- ultrashort X-ray pulse generation
- more detailed numerical modeling
- experimental implementation
- these are only the first examples (orientation of molecules, accelerating molecules, or nanostructures, etc.)



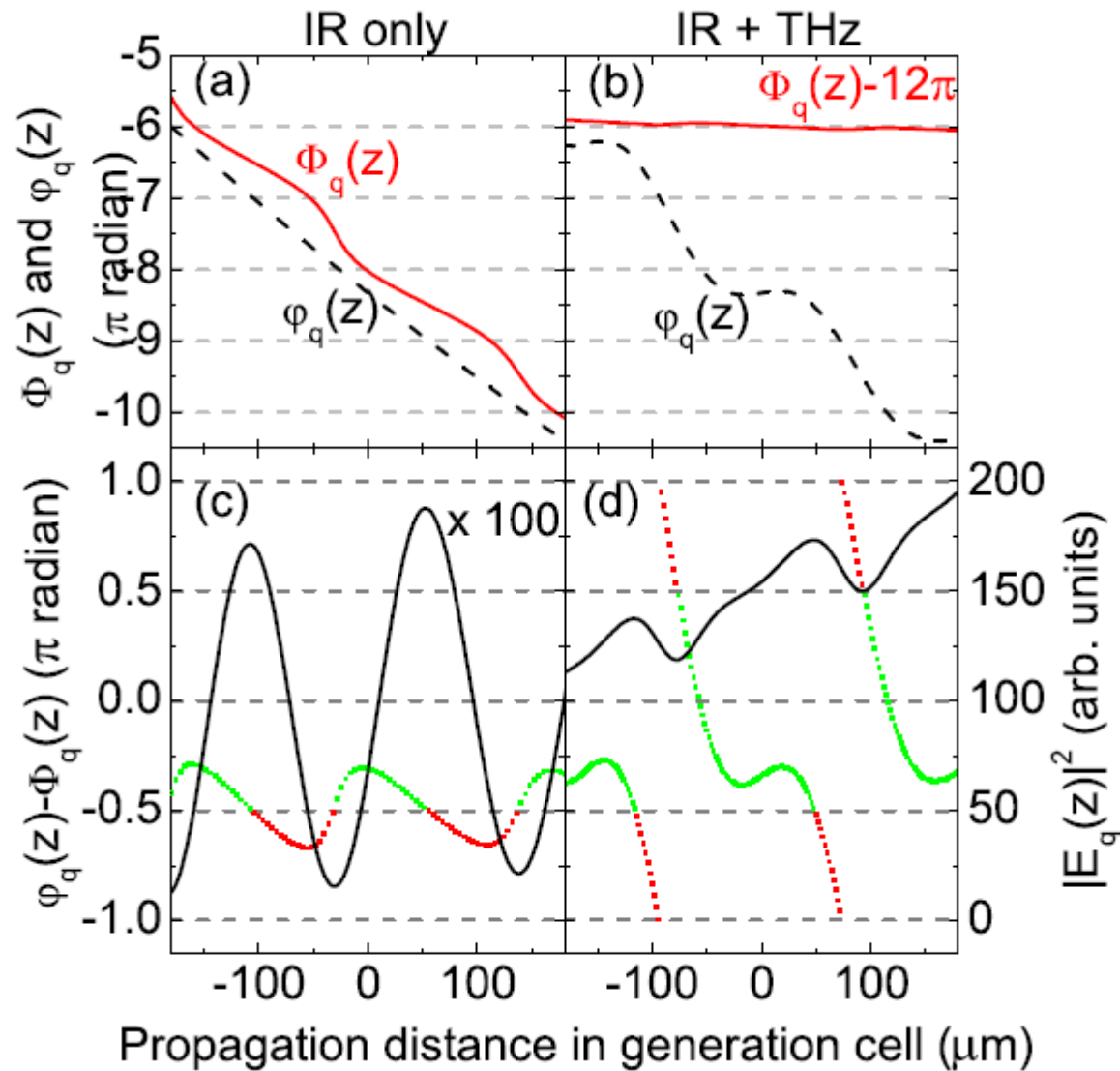
Acknowledgement: Hungarian Scientific Research Fund (OTKA), grant number 76101, 78262, 101846, SROP-4.2.1.B-10/2/KONV-2010-0002, and hELIos ELI_09-01-2010-0013

Single-cycle EUV pulse generation by coherent Thomson-scattering

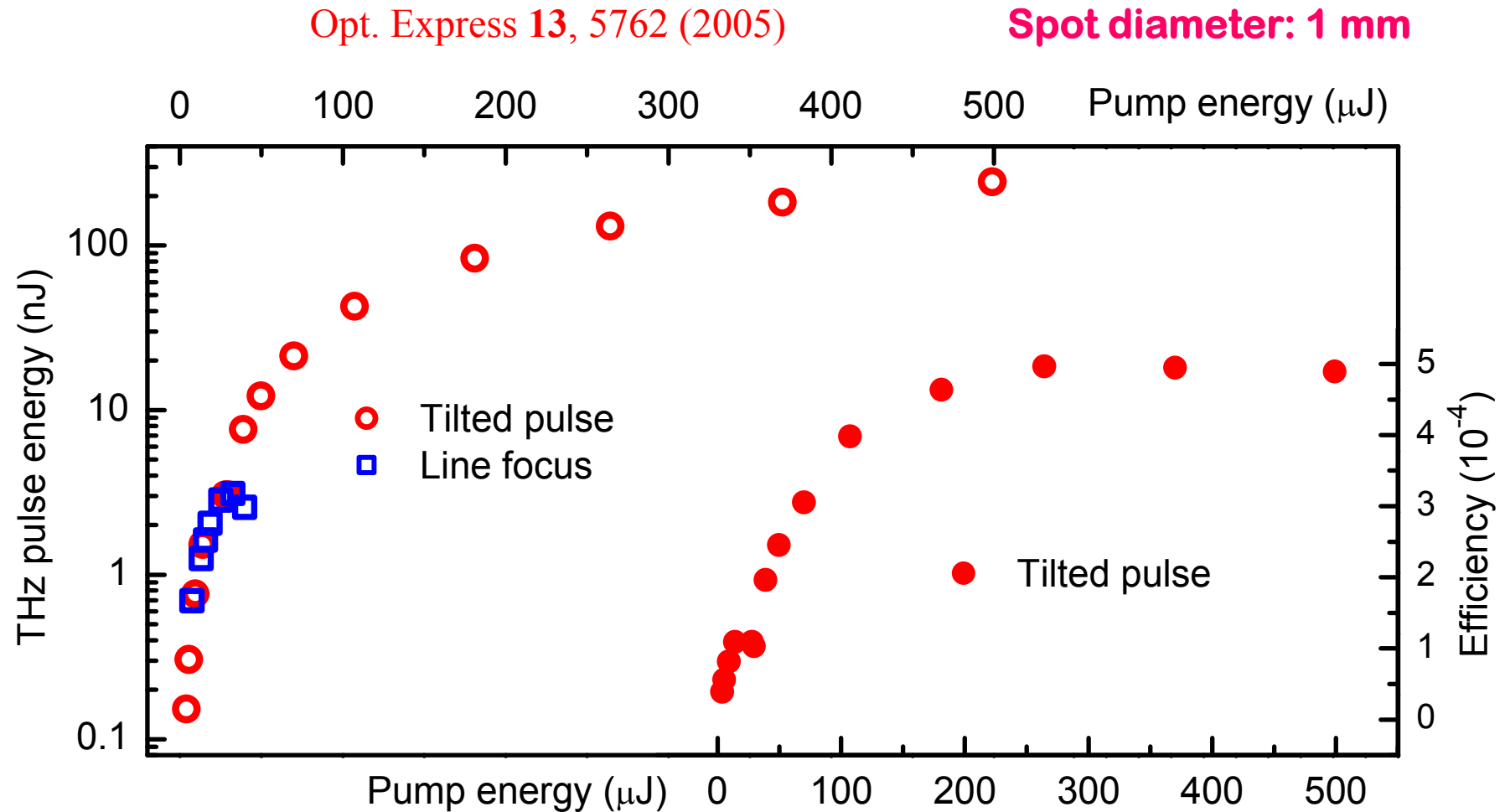


Quasi-phase-matched attosecond pulse generation

Spatial modulation by (chirped) THz pulse



Scaling-up the THz pulse energy

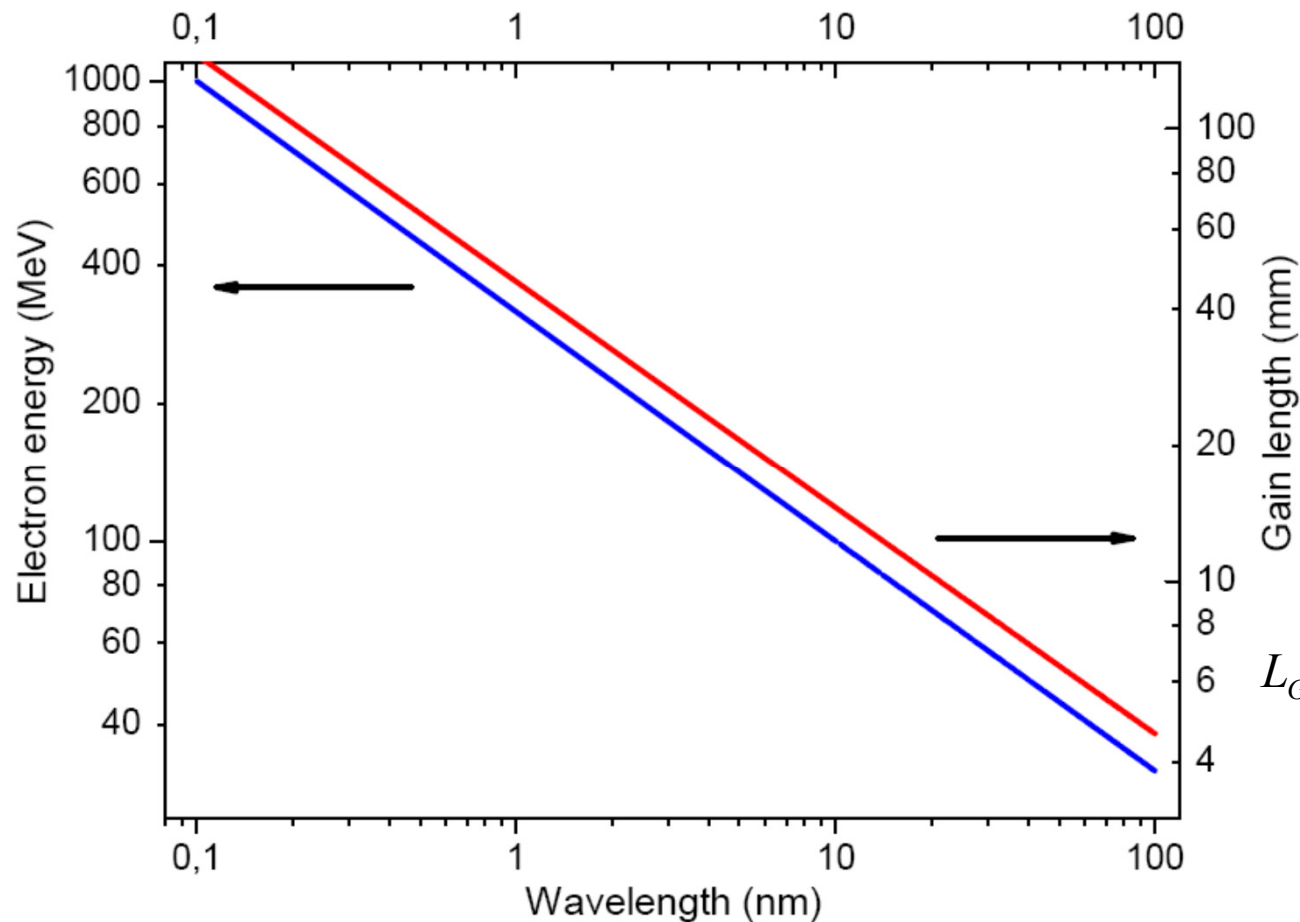


$E_{\text{THz}} = 240 \text{ nJ}$, $\eta_{\text{THz}} = 10 \%$, electric field (unfocused) 0.1 MV/cm

THz Workshop, Argonne National Laboratory, July 30-31, 2012

THz driven undulator for UV – X-ray generation

$\lambda_u = 600 \mu\text{m}$, $E_{\text{THZ}} = 40 \text{ MV/cm}$, $K = 0.747$, $Q = 420 \text{ pC}$, $\sigma_L = 7 \mu\text{m}$, $\sigma_T = 30 \mu\text{m}$



$$\lambda_r = \frac{\lambda_u}{2\gamma^2} \left[1 + \frac{K^2}{2} \right]$$

$$K = \frac{\lambda_u \cdot e \cdot E_{\text{THZ}}}{2\pi \cdot mc^2}$$

$$L_G \propto \lambda_u^{5/6}$$

$$L_G = \frac{\sqrt[3]{\lambda_u} \cdot \gamma \cdot \sigma_T^{2/3} \cdot (1 + \Lambda)}{\sqrt[3]{\pi} \cdot \sqrt{3} \cdot K^{2/3}} \cdot \sqrt[3]{\frac{I_A}{I}}$$

$$\frac{\Delta\gamma}{\gamma} \propto \sqrt[6]{\lambda_u} \quad \varepsilon_T \propto \lambda^{3/2}$$



J. Hebling et al.: arXiv 1109.6852 (2011)

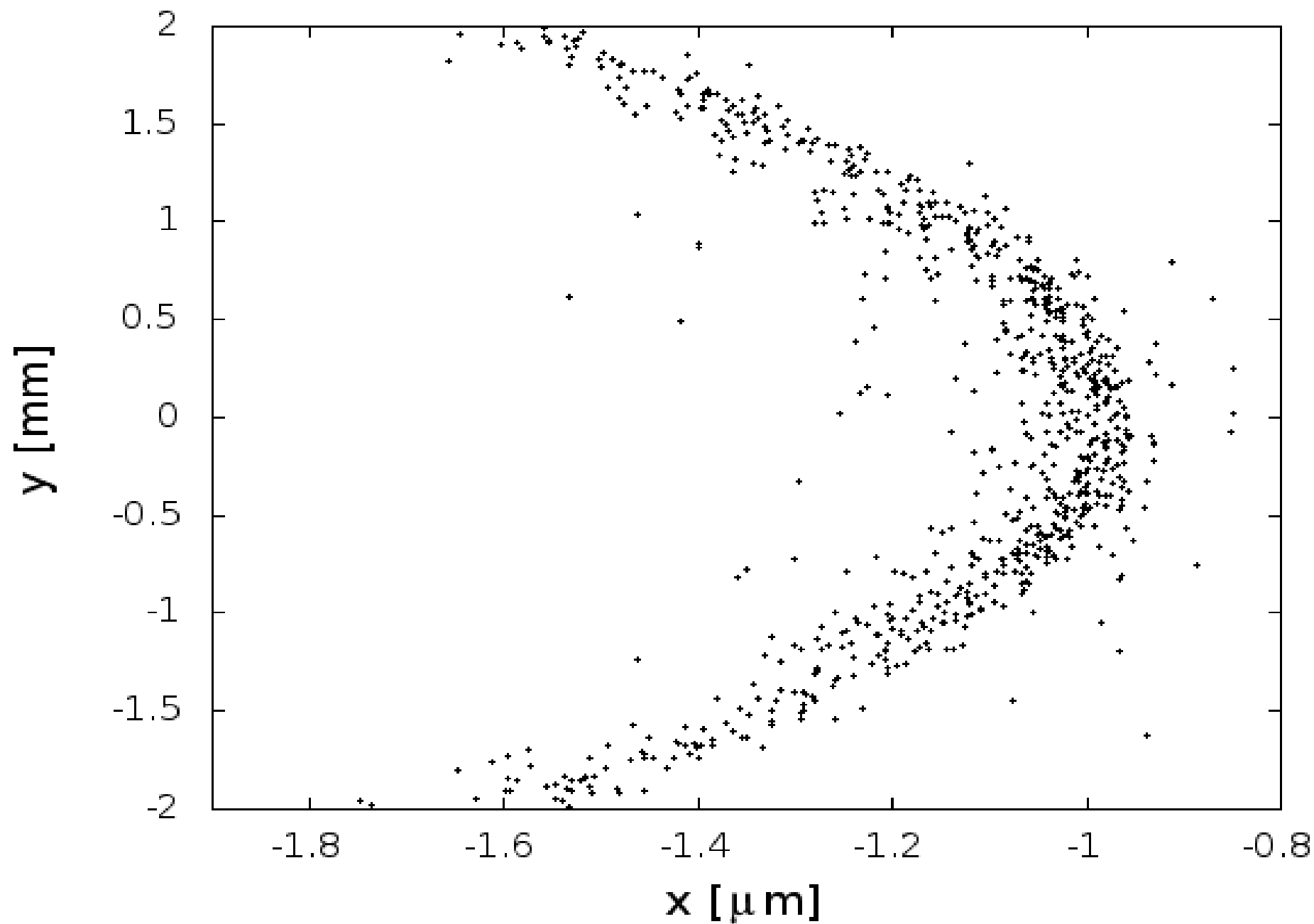
$W_{\text{THz}} = 14 \text{ J} !!$

ELI-ALPS light sources

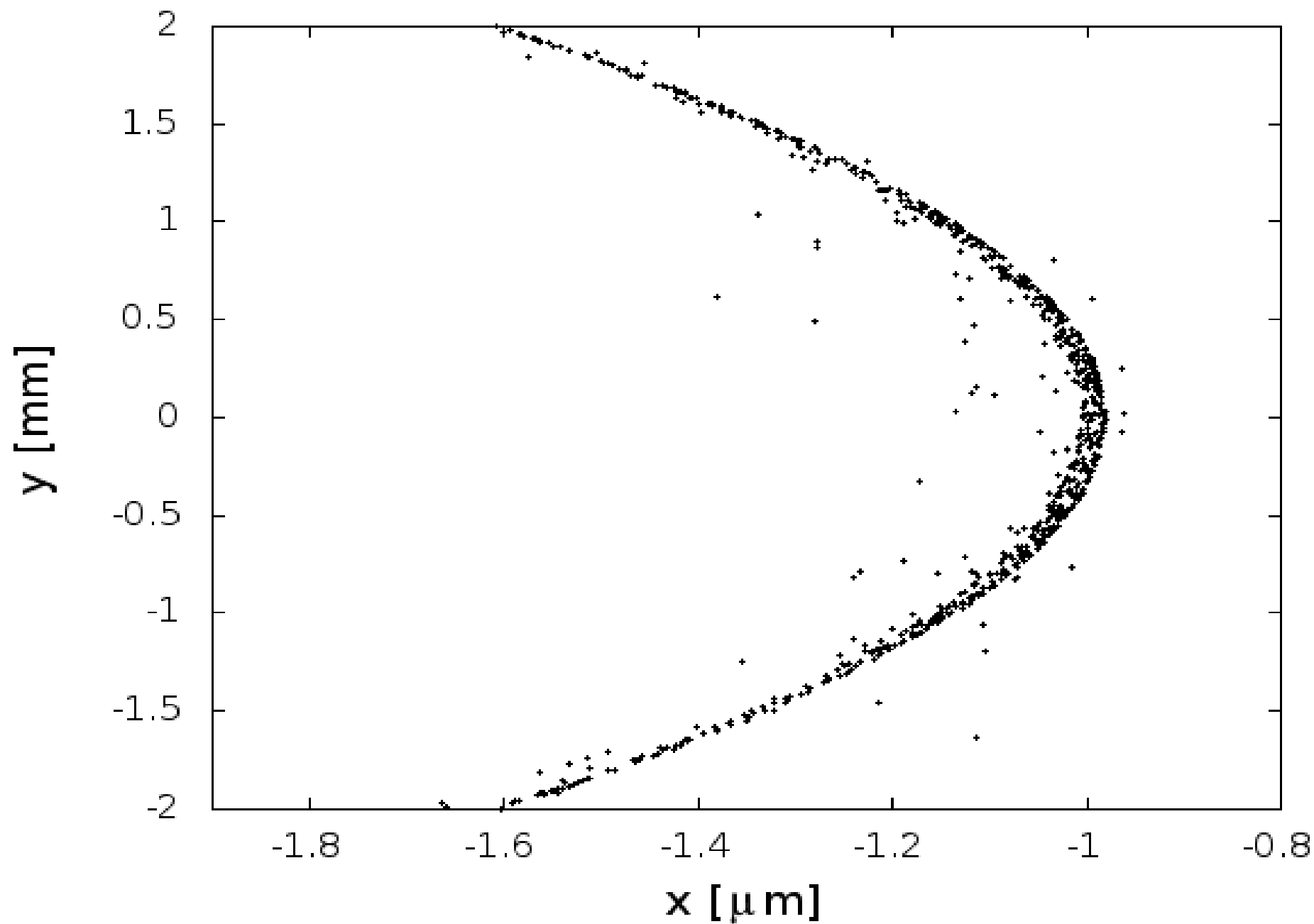
Primary sources	Peak /average power	Repetition rate	Pulse Energy	Pulse Duration	Spectral Range	Secondary sources
ALPS-WB	2-3TW / 300-500W	100kHz	3-5mJ	1.5-3fs	0.3-1.3μm	-T0 FIR/THz 0.3-3THz, 100μm-1mm, 1.24-12.4meV / 0.1-1 mJ, 10-30 MV/cm -T1 MIR 3-30THz, 10-100μm, 12.4-124meV / 10-100μJ, 3-300 MV/cm -T2 NIR 30-300THz, 1-10μm, 0.12-1.24eV / >100μJ, 0.3-30 GV/cm; as an alternative an IR OPCPA system can provide similar parameters -P1 UV/XUV 1-25PHz, 12-300nm 4-100eV / 100-10nJ -P2 SXR 25-100PHz, 3-12nm, 100-400eV / >1nJ
SYLOS (<i>Single-cycle OPCPA System</i>)	100TW / 300-1000W	1kHz	0.3-1J	3-5fs	0.5-1.3μm	-S1 10-100eV, 12-120nm / 100-1μJ -S2 100-1000eV, 1.2-12nm / 1-0.1μJ -H1 1-10keV, 0.12-1.2nm / 100-1nJ -H2 10-100keV, 0.12-1.2A / >1nJ
ALPS-HF	1-3PW / 10-100W (>2PW / >400W)	1-10Hz (10Hz)	3-10J (>40J)	3-5fs (<20fs)	0.5-1μm / 0.7-1.3μm (0.7-0.9μm)	Controlled electric field: (3-5)x10 ¹⁴ V/m Normalized vector potential: 60-110/80-140 Peak intensity: (1-3)x10 ²² W/cm ² Ponderomotive potential: (0.6-2)x10 ²² / (1-3)x10 ²² W/cm ² × μm ²



$\dot{u}t=3.238\text{m}$, $\gamma_{\min}=19.93$, $\gamma_{\max}=20.059$



$\dot{u}t=3.258m$, $\gamma_{\min}=19.93$, $\gamma_{\max}=20.059$



Optical rectification

Conversion efficiency depends on:

- ω_{THz} ($\eta \sim \omega_{\text{THz}}^2$)

- material parameters

- phase-matching \rightarrow velocity matching: $v_{\text{vis}}^{\text{gr}} = v_{\text{THz}}^{\text{ph}}$

$$\eta_{\text{THz}} = \frac{2\omega^2 d_{\text{eff}}^2 L^2 I}{\epsilon_0 n_v^2 n_{\text{THz}} c^3} \cdot \exp\left[-\alpha_{\text{THz}} L/2\right] \cdot \frac{\sinh^2\left[\alpha_{\text{THz}} L/4\right]}{\left[\alpha_{\text{THz}} L/4\right]^2}$$

$$\alpha_{\text{THz}} L \ll 1$$

$$\eta_{\text{THz}} = \frac{2\omega^2 d_{\text{eff}}^2 L^2 I}{\epsilon_0 n_v^2 n_{\text{THz}} c^3}$$

$$FOM_{\text{NA}} = \frac{d_{\text{eff}}^2 L^2}{n_v^2 n_{\text{THz}}}$$



$$\alpha_{\text{THz}} L \gg 1$$

$$\eta_{\text{THz}} = \frac{8\omega^2 d_{\text{eff}}^2 I}{\epsilon_0 n_v^2 n_{\text{THz}} \alpha_{\text{THz}}^2 c^3}$$

$$FOM_A = \frac{4d_{\text{eff}}^2}{n_v^2 n_{\text{THz}} \alpha_{\text{THz}}^2}$$

Figure of merits (FOMs) supposing 2 mm long crystals

Material	d_{eff} (pm/V)	n_{800nm}^{gr}	n_{THz}	$n_{1.55\mu m}^{gr}$	α_{THz} (cm ⁻¹)	FOM (pm ² cm ² /V ²)
CdTe	81.8		3.24	2.81	4.8	11.0
GaAs	65.6	4.18	3.59	3.56	0.5	4.21
GaP	24.8	3.67	3.34	3.16	0.2	0.72
ZnTe	68.5	3.13	3.17	2.81	1.3	7.27
GaSe	28.0	3.13	3.27	2.82	0.5	1.18
sLiNbO ₃ sLN 100K	168	2.25	4.96	2.18	17 4.8	18.2 48.6
DAST	615	3.39	2.58	2.25	50	41.5



Velocity matching condition:

$$v_{NIR}^{gr} = v_{THz}^{ph} \Rightarrow n_{NIR}^{gr} = n_{THz}$$

THz Workshop, Argonne National
Laboratory, July 30-31, 2012

Tilted-Pulse-Front Pumping (TPFP) Setup

Fülöp & Hebling,
in: Recent Optical and
Photonic Technologies, 2010

$$\tan \gamma = -\frac{n}{n_g} \lambda \frac{d\varepsilon}{d\lambda}$$

↑
pulse front tilt

↑
angular dispersion

